"Quantum Gravity corrections at the Planck time"

Elias C. Vagenas



Work in collaboration with : Spyros Basilakos, RCAAM, Academy of Athens Saurya Das, Univ. of Lethbridge, Canada



Related references:

- S. Das and ECV, Phys.Rev.Lett.101:221301,2008. e-Print: arXiv:0810.5333 [hep-th]
- S. Das and ECV, Can.J.Phys.87:233-240,2009. e-Print: arXiv:0901.1768 [hep-th]
- A.F. Ali, S. Das and ECV, Phys.Lett.B678:497-499,2009.
 e-Print: arXiv:0906.5396 [hep-th]
- A.F. Ali, S. Das and ECV, to appear in Phys.Lett.B 2010. e-Print: arXiv:1005.3368 [hep-th]



<u>Outline</u>

- Motivation
- Heisenberg Uncertainty Principle and Planck era
- Generalized Uncertainty Principle (GUP)
- GUP at the Planck time
- Conclusions



Motivation

Gravity is a universal and fundamental force. Anything that has energy creates gravity and is affected by it.

Newton's constant G often means that associated classical effects are too weak to be measurable.

Prediction of candidate theories of QG (such as String Theory) and black hole physics : *minimum measurable length.*

Model-independent prediction and can be understood as follows: HUP breaks down for energies close to Planck scale (Schwarzschild radius is comparable to the Compton wavelength).



So at Planck scales the suggestion is that GUP holds:

$$\Delta x \ge \frac{\hbar}{\Delta p} + \alpha' \frac{\Delta p}{\hbar}$$

(E. Witten, Physics Today, 1996)

Recently it was shown that certain effects of QG are also universal and can affect almost any system with a well-defined Hamiltonian (quantum effects are generically small to be observed in experiments but future experiments could make them measurable).

(A.F. Ali, S. Das, ECV, Phys. Lett. B 678, 497 (2009))

The scope is to investigate in a cosmological setup if these GUPcorrections are assigned to physical quantities at Planck time.

Heisenberg Uncertainty Principle and Planck era

$\Delta x \Delta p \geq \hbar$

In order to derive the energy-time uncertainty principle, one employs the relations:

$$\Delta x \approx c \, \tau, \quad \Delta p \approx \frac{\Delta E}{c}$$

and thus the energy -time uncertainty principle reads:

$$\Delta E \tau \ge \hbar$$



At Planck era one can easily derive the energy-time uncertainty principle as follows:

$$E_{Pl} = M_{Pl}c^2, M_{Pl} = \rho_{Pl}l_{Pl}^3$$

$$\mathcal{O}_{Pl} \approx \frac{1}{Gt_{Pl}^2}$$

1

(1st Friedmann eqn)

Then set

$$\Delta E \approx E_{Pl} \qquad \tau \approx t_{Pl}$$



 $\rho_{Pl} l_{Pl}^{3} c^2 t_{Pl} \ge \hbar$ $\Rightarrow \frac{l_{Pl}^{3}}{Gt_{Pl}^{2}}c^{2}t_{Pl} \ge \hbar$

$$\Rightarrow \frac{c^5}{G} t_{Pl}^{2} \ge \hbar$$



(P. Coles, F. Lucchin, "Cosmology: The Origin and the Evolution of the Cosmic Structure", 1995)



Generalized Uncertainty Principle (GUP)

Recently it was proposed (by Ahmed F Ali, Saurya Das and ECV) a new for of commutators that are consistent with String Theory, Black Hole Physics and DSR:

$$[x_i, p_j] = i\hbar \left(\delta_{ij} - (p\delta_{ij} + \frac{p_i p_j}{p}) + \alpha^2 (p^2 \delta_{ij} + 3p_i p_j)\right)$$

$$\Delta x \Delta p \ge \frac{\hbar}{2} \left(1 - 2\alpha \langle p \rangle + 4\alpha^2 \langle p^2 \rangle \right)$$

$$\Delta x \ge (\Delta x)_{\min} \approx \alpha_0 l_{Pl}$$
$$\Delta p \le (\Delta p)_{\max} \approx \frac{M_{Pl}c}{\alpha_0}$$
where $\alpha = \frac{\alpha_0 l_{Pl}}{\hbar}$



 $\Delta x \Delta p \ge \frac{\hbar}{2} \left(1 - 2\alpha \langle p \rangle + 4\alpha^2 \langle p^2 \rangle \right)$ $\oint \Delta E \tilde{t}_{Pl} \ge \frac{\hbar}{2} \left(1 - 2\alpha \frac{\Delta E}{c} + 4\alpha^2 \frac{(\Delta E)^2}{c^2} \right)$



$\Delta E = \widetilde{E}_{Pl} = \widetilde{M}_{Pl} c^2, \ \widetilde{M}_{Pl} = \widetilde{\rho}_{Pl} \widetilde{l}_{Pl}^3$ $\widetilde{\rho}_{Pl} \approx \frac{1}{G\widetilde{t}_{Pl}^2}$

Finally one gets:

 $\tilde{t}_{Pl} = t_{Pl} f_+(\alpha_0)$



$$f_{\pm}(\alpha_{0}) = \left[\frac{(1+\alpha_{0}) \pm \sqrt{(1-\alpha_{0})(1+3\alpha_{0})}}{4\alpha_{0}^{2}}\right]^{1/2}$$
$$f_{\pm}(1) = \frac{\sqrt{2}}{2} , \quad f_{\pm}(-\frac{1}{3}) = \sqrt{\frac{3}{2}}$$
and
$$\lim_{\alpha_{0} \to 0} f_{-}(\alpha_{0}) = \frac{\sqrt{2}}{2}$$

SO $f_{-}(\alpha_{0})$ does not pay any significant role since

$$\widetilde{t}_{Pl} \approx t_{Pl}$$



On the contrary, for the case of $~f_+(lpha_0)$

and for small values of α_0

$$f_+(\alpha_0) \approx \frac{1}{2\alpha_0}$$

As an example

NEB14

8-11 June 2011, Ioannina,



goes rapidly to infinity!!

$$\alpha_{0} = O(10^{-2} - 10^{-4})$$
we find that $f_{+} \approx \frac{\tilde{t}_{Pl}}{t_{Pl}} \approx 10^{2} - 10^{4}$!!!
$$So \quad \frac{\tilde{\rho}_{Pl}}{\rho_{DE}} \approx O(10^{115} - 10^{119})$$

$$\tilde{\sigma}_{Pl} \approx \sigma_{Pl}$$
 !!!!

Conclusions

 Under specific circumstances the main Planck quantities in the context of GUP are LARGER than those defined in the context of HUP by a factor of

$$f_+ \approx 10^2 - 10^4$$

• Dimensionless entropy enclosed in the cosmological horizon does not "feel" the GUP effects and thus information remains the same.

•We expect modifications in the framework of Cosmology since changes in the Planck epoch will be inherited to late universe through QG.

