Constraints on sub-eV physics beyond the SM from cosmological distance measurements

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Outline



Sub-eV Physics beyond the SM: ALPs & MCPs in Field & String Theory



Cosmological Distance Measures







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The QCD Axion

QCD allows for a CP-violating term:

$$\mathcal{L}_{CP} = \frac{\alpha_s}{4\pi} \,\theta \,\mathrm{tr}G_{\mu\nu}\tilde{G}^{\mu\nu}$$

Parameter θ constrained experimentally:

$$\left|\bar{\theta}\right| \lesssim 10^{-10}$$

unnaturally small

coupling to

<u>Peccei-Quinn</u>: Promote θ to a dynamical field, the axion a, with shift symmetry $a \rightarrow a + \text{const}$:

$$\mathcal{L}_{a} = \frac{1}{2} \partial_{\mu} a \partial^{\mu} a + \frac{\alpha_{s}}{4\pi f_{a}} a \operatorname{tr} G_{\mu\nu} \tilde{G}^{\mu\nu} + \frac{\alpha}{8\pi f_{a}} a F_{\mu\nu} \tilde{F}^{\mu\nu} + \mathcal{L}_{\mathrm{int}} [\partial_{\mu} a / f_{a}, \psi]$$

Non-trivial potential around $\langle a \rangle = 0$, axion is a PNG boson with parametrically small mass:

$$m_a \simeq 0.6 \,\mathrm{meV} \times \left(\frac{10^{10} \mathrm{GeV}}{f_a}\right)$$

Axions in String Theory

Axion-Like Particles (ALPs) arise in String Theory as 0-modes of antisymmetric tensor fields

• <u>Type II</u>: bosonic action for a Dp-brane has two contributions

$$S_{p} = -T_{p} \left(\int d^{p+1} \xi \, e^{-\phi} \sqrt{\det(g + B + 2\pi\alpha' F)} + i \int \sum_{q} C_{q} \wedge e^{B + 2\pi\alpha' F} \right)$$

DBI piece: includes $F_{\mu\nu}F^{\mu\nu}$ *WZ piece: includes* $aF_{\mu\nu}\tilde{F}^{\mu\nu}$

Axion decay const set by the string scale:

$$f_a \sim \frac{M_s}{g_s} \sim 10^{4-17} \,\mathrm{GeV}$$

• <u>Heterotic</u>:

$$S_{H} = \frac{2\pi M_{s}^{8}}{g_{s}^{2}} \int d^{10}x \sqrt{-g}R - \frac{M_{s}^{6}}{2\pi g_{s}^{2}} \int \frac{1}{4} \mathrm{tr}F \wedge \star F - \frac{M_{s}^{4}}{g_{s}^{2}} \int \frac{1}{2} \mathrm{tr}H \wedge \star H + \dots$$

with axion decay const of order $f_a = \frac{\alpha_{YM}M_P}{2\pi\sqrt{2}} \sim 10^{16} \,\mathrm{GeV}$

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 $\tilde{r}\mu\nu$

Other sub-eV particles

• <u>Hidden U(1)</u> fields, coupling to photons via:

$$\mathcal{L}_{\rm mix} = \frac{\chi}{2eg_h} F_{\mu\nu} X^{\mu\nu}$$

Mixing parameter $10^{-17} \lesssim \chi \lesssim 10^{-5}$

- <u>Hidden matter</u> charged under theses U(I)s Small (unquantised) charge $\epsilon \sim \chi g_h$ aka Mini-Charged Particles (MCPs)
- <u>Chameleons</u> can couple to (ordinary) photons

Constraints on these particles come predominantly from low-E experiments (lasers, microwave cavities, torsion balances,...)

Here we will constrain ALPs (& MCPs) by comparing cosmological distance measurements

Distance Measures in Cosmology

<u>Luminosity distance</u>:

$$d_L(z) = (1+z)\frac{c}{H_0} \int_0^z dz' \left[\Omega_m (1+z)^3 + \Omega_V (1+z)^{3(1+w)}\right]^{-1/2}$$

Inferred from standard candles, notably la SNae

 <u>Ang. diameter distance</u> related through Etherington relation: (from standard rulers)

$$d_L(z) \neq (1+z)^2 d_A(z)$$

If photon number conservation is violated, there will be a mismatch in the above due to a non-trivial "opacity" τ :

$$d_{L,obs}(z) = d_{L,true}(z) e^{\tau(z)}$$

This can happen if photons are converted to ALPs along line of sight

Constraining opacity & ALPs



Any ALP coupling to photons via will produce non-trivial opacity.

 $\frac{1}{4M}F_{\mu\nu}F^{\mu\nu}\phi \quad \text{or} \quad \frac{1}{8M}\epsilon_{\mu\nu\kappa\lambda}F^{\mu\nu}F^{\kappa\lambda}\phi$

Can constrain jointly ALP coupling and cosmological parameters by using SN and H(z) (or BAO) data.

Method

Run likelihood analysis for flat ACDM models in (τ, Ω_m, H_0)

Constrain opacity parameter(s) by marginalising over cosmologies:

$$P(\tau|S,E) = \int_{\Omega_m} \int_{H_0} P(\tau,\Omega_m,H_0|S) P(\Omega_m,H_0|E) d\Omega_m dH_0$$
$$dL_{obs}(z) = (1+z)^2 e^{\tau} d_A(z)$$

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$$(d_{L,obs}(z)) = (1+z)^2 e^{\tau} (d_A(z))$$

Method

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$$dL_{obs}(z) = (1+z)^2 e^{\tau} d_A(z)$$

•For ALPs:

$$e^{-\tau} \equiv \mathcal{P}(z) = A + (1 - A) \exp\left(-\frac{P}{H_0 L} \frac{H(z) - H_0}{\Omega_m H_0}\right)$$
Initial SN flux mix: $A = \frac{2}{3}\left(1 + \frac{I_{\phi}}{I_{\gamma}}\right)$
Photon-axion conversion probability

•For MCPs:
 $e^{-\tau} \equiv \mathcal{P}(z) = \exp\left(-\kappa y(z)\right)$
Rate of $\gamma \to \psi \bar{\psi}$

Axion-Like Particles (incl. Chameleons)

Simplest ALPs (A=2/3)

If ALPs have no other interactions than 2-photon couplings, then they do NOT contribute to SN luminosity (i.e. $A \approx 2/3$)

This constraint on P translates into a constraint on the axion-photon coupling scale M, intergalactic magnetic field B and the electron density

Constraining the Axion Coupling

The photon-axion conversion probability P, depends on the combination B/M and on the electron number density:

Bounds from cosmological distances already competitive with CMB and Quasar bounds

Stronger constraints on axion-photon coupling will be placed as intergalactic magnetic fields are better understood and observational data improve

Mini-Charged Particles

For MCPs ψ the rate κ of the $\gamma \to \psi \overline{\psi}$ process depends on the MCP charge and mass, as well as the magnetic field:

- Constraints from cosmology competitive with lab experiments
- H(z) significantly improves constraints
- Expect dramatic improvement with future H(z)/BAO data

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Forecasts (BAO & SN)

Dramatic improvement on these constraints expected with future BAO (notably EUCLID) and SN (SNAP) missions

Summary

- Vast quantity of high quality cosmo data fast approaching: CMB, BAOs, Gravitational waves, 21 cm,...
- Fruitful interplay between HEP/cosmo theory and cosmological observation (cf compactification scales from inflation!)
- New physics at sub-eV scales (notably ALPs & MCPs) generic in fundamental theory
- Cosmological distance measurements are already placing bounds on microphysical (& astrophysical) parameters
- Dramatic improvement expected as new data arrives and astrophysics better understood